

Problem Sheet 5
Quantum channels

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(*Hint:* While the first two excersises are more abstract the last one discusses simple examples and consists of hands-on calculations. If you get stuck somewhere on this sheet, try to jump ahead to the examples or vice versa.)

Let \mathcal{X} and \mathcal{Y} be two Hilbert spaces and $L(\mathcal{X})$ and $L(\mathcal{Y})$ denote the linear operators on the Hilbert spaces. In the lecture, you got to know quantum channels as those linear maps $T : L(\mathcal{X}) \rightarrow L(\mathcal{Y})$ such that the map $T \otimes \mathbb{1}_d$ maps quantum states to quantum states where $\mathbb{1}_d$ acts on an additional Hilbert space of arbitrary dimension d . You have seen that this amounts to requiring that quantum channels are linear maps that are completely positive and trace preserving (CPT). If not specified otherwise, on this sheet, T will always denote such a quantum channel.

You also got to know various representations of quantum channels

- The *Kraus representation*. A map $T \in L(L(\mathcal{X}), L(\mathcal{Y}))$ is CPT *iff* there exist a set of linear operators $\{K_i\}_i$ with $K_i \in L(\mathcal{X}, \mathcal{Y})$ fulfilling $\sum_i K_i^\dagger K_i = \mathbb{1}$ such that

$$T(X) = \sum_i K_i X K_i^\dagger. \quad (1)$$

- The *Stinespring representation*. There exists an isometry $V \in L(\mathcal{X}, \mathcal{Y} \otimes \mathcal{Z})$ or, equivalently¹, an arbitrary reference state $|0\rangle \in \mathcal{Z}'$ and a corresponding unitary operator $U \in U(\mathcal{X} \otimes \mathcal{Z}')$ with $\mathcal{Y} \otimes \mathcal{Z} \cong \mathcal{X} \otimes \mathcal{Z}'$ such that

$$T(X) = \text{Tr}_{\mathcal{Z}}[V X V^\dagger] = \text{Tr}_{\mathcal{Z}}[U(X \otimes |0\rangle\langle 0|)U^\dagger]. \quad (2)$$

- The *Choi-Jamiołkowski representation*. $J(T) \in \mathcal{Y} \otimes \mathcal{X}$

$$J(T) := (T \otimes \mathbb{1}) |\omega\rangle\langle\omega|, \quad (3)$$

where $|\omega\rangle = \frac{1}{\sqrt{d}} \sum_i |ii\rangle$ is the maximally entangled state.

In this problem sheet, we will show the equivalence between those representations explicitly and consider some examples.

1. On the Kraus representation of quantum channels

The operational meaning of Kraus operators can be understood in the following setting in which, for simplicity, we restrict ourselves to quantum channels with the same input and output space $L(\mathcal{X})$. Suppose we apply a unitary U to the joint system and environment in the state $\rho \otimes |0\rangle\langle 0| \in L(\mathcal{X} \otimes \mathcal{Z})$, where $|0\rangle \in \mathcal{Z}$ is some reference state, and then we measure system \mathcal{Z} in the computational basis.

- a) Show that the action of the unitary on the joint system can be written as

$$U(\rho \otimes |0\rangle\langle 0|)U^\dagger = \sum_{kl} E_k \rho E_l^\dagger \otimes |k\rangle\langle l|,$$

with respect to the basis $\{|i\rangle\}_i$ on the second system.

¹Here, we use that any isometry $V : \mathcal{X} \rightarrow \mathcal{X} \otimes \mathcal{Z}'$ can be written as $A = U(\mathbb{1} \otimes |0\rangle)$ with an arbitrary reference state $|0\rangle$ and a corresponding unitary $U \in U(\mathcal{X} \otimes \mathcal{Z}')$.

- b) Now, we perform a von-Neumann measurement on \mathcal{Z} in the same basis. Determine the post-measurement state conditioned on outcome i . What is the probability of obtaining outcome i ?
- c) Give the corresponding operational interpretation of the Kraus operators E_k and the unitary U .
- d) Now, suppose we want to implement a von-Neumann measurement on \mathcal{X} in this manner. Characterize the unitaries $U \in U(\mathcal{X} \otimes \mathcal{Z})$ on the joint system that give rise to this situation. Give an example for the case of two qubits.
- e) What do the operators E_k have to satisfy such that one can reverse the channel after having (destructively) measured outcome k on \mathcal{Z} ?

Finally, we will show some properties of the Kraus representation

- f) Let $\{K_i\}_{i=1}^N$ and $\{\tilde{K}_j\}_{j=1}^N$ be two sets of linear operators in $L(\mathcal{X}, \mathcal{Z})$ fulfilling the completeness relation of Kraus operators. Show that if the two sets are related by a unitary transformations $U \in U(N)$ such that $\tilde{K}_i = \sum_j U_{ij} K_j$, the channels represented by the sets coincide.
- g) Show that all equal-sized Kraus representations of a given channel T are related via a unitary transformation.

Hint: Relate the Kraus representation two low-rank matrix factorisations of the Choi matrix.

2. Equivalence between representations of quantum channels

Let us first show that the Choi-Jamiołkowski map $J : L(L(\mathcal{X}), L(\mathcal{Y})) \rightarrow L(\mathcal{Y} \otimes \mathcal{X})$ is a linear bijection between the CPT maps on the one hand and the set of quantum states on $\mathcal{Y} \otimes \mathcal{X}$ on the other hand.

- a) Show that the inverse map can be defined by $\tilde{T}(X) = \text{Tr}_{\mathcal{X}}[J(T)(\mathbb{1}_{\mathcal{Y}} \otimes X^T)]$ and makes J a bijection.

Let $\rho_T \in \mathcal{Y} \otimes \mathcal{X}$ be the Choi-Jamiołkowski state corresponding to the quantum channel T .

- b) Determine a set of Kraus operators representing T .
- c) Determine a unitary U_T representing T via the Stinespring representation.

Now, let U_T be a unitary representing T in the Stinespring representation.

- d) Determine the Choi-Jamiołkowski state representing T .

The rank of a quantum channel is defined as the rank of its Choi matrix.

- e) Show that a quantum channel with rank r can be represented as a Stinespring dilation using an auxiliary system of dimension r .

3. Examples of quantum channels

Now we are ready to look at some examples. The following maps are important so-called noise channels

$$F_{\epsilon}(A) := \epsilon X A X + (1 - \epsilon) A$$

$$D_{\epsilon}(A) := \epsilon \text{Tr}[A] \frac{\mathbb{1}}{d} + (1 - \epsilon) A$$

$$A_{\epsilon}(A) := \epsilon \text{Tr}[A] |0\rangle\langle 0| + (1 - \epsilon) A,$$

where $\epsilon \in [0, 1]$.

- a) For each channel, show that it is CPT.

- b) For each channel, give its Choi-Jamiołkowski state, a Kraus representation and a Stinespring representation.

Hint: It may help to consider $\epsilon = 1$ in a first step and then generalize to arbitrary $\epsilon \in [0, 1]$.

- c) Give a physical interpretation and a good name for each channel.