#### Freie Universität Berlin

## **Tutorials on Quantum Information Theory**

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## Problem Sheet 6 Witnessing entanglement

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## 1. Warm-up: Partial transpose

In the entanglement theory of bi-partite systems the partial transpose criterion plays a prominet rôle. Let  $T: L(\mathcal{H}) \to L(\mathcal{H})$  be the transposition map  $X \mapsto T(X) = X^T$ . The partial transpose is the map  $T: L(\mathcal{H} \otimes \mathcal{H}) \to L(\mathcal{H} \otimes \mathcal{H})$ . Let  $(\cdot, \cdot)$  be the Hilbert-Schmidt inner-product on  $L(\mathcal{H})$  defined as  $(X, Y) = \text{Tr}(X^TY)$ . The adjoint  $\Lambda^*$  of a map  $\Lambda: L(\mathcal{H}) \to L(\mathcal{H})$  with respect to  $(\cdot, \cdot)$  is defined such that  $(\Lambda(X), Y) = (X, \Lambda^*(Y))$  holds for all  $X, Y \in L(\mathcal{H})$ .

a) Show that  $\mathbb{1} \otimes T$  is self-adjoint, i.e.  $(\mathbb{1} \otimes T)^* = \mathbb{1} \otimes T$ .

# 2. An entanglement witness for the maximally entangled state

Let  $|\Omega\rangle$  be the maximally entangled state.

- a) Show that  $W = 1 d |\Omega\rangle\langle\Omega|$  is a witness for  $|\Omega\rangle$ .
- b) Give an example of an entangled state that is not detected by W.

### 3. The reduction map as a witnesss

The reduction map is defined as  $\Lambda_R(X) = \text{Tr}(X)\mathbb{1} - X$ .

- a) Show that  $\Lambda_R$  is positive but not completely positive, in other words it is a witness.
- b) Give at least one example for states that are detectable by  $\Lambda_R$ .
- c) Give at least one example for entangled states that are not detected by  $\Lambda_R$ .
- d) What is the observable witness associated to the positive map by the Choi-Jamiolkowski-isomorphism?

A map  $\Lambda$  is called *decomposable* if it can be written as  $\Lambda = P_1 + P_2 \circ T$ , where  $P_1, P_2$  are completely positive maps and T is the transpose.

- e) Show that any state that is detected by  $\Lambda_R$  can also be detected by the partial transpose criterion.
  - *Hint:* Argue first that  $\Lambda_R$  is decomposable.
- f) Translate the condition of a map  $\Lambda$  being decomposable to a criterion for the observable witness  $J(\Lambda^*)$ . What is the implication of a self-adjoint observable witness being decomposable in this sense?

### 4. Constructing entanglement witnesses from the partial transpose

In the lecture, we saw that for any positive but not completely positive map  $\Lambda$ ,  $J(\Lambda^*)$  is also witness. But  $J(\Lambda^*)$  is not necessarily detecting all the states that  $\Lambda$  detects. Let  $\Lambda$  detect  $\rho_e$  and let  $|\eta\rangle$  be an eigenvector with negative eigenvalue of  $(\mathbb{1} \otimes \Lambda)(\rho_e)$ .  $|\Omega\rangle$  shall be the maximally entangled state.

a) Show that if  $J(\Lambda^*)$  detects  $\hat{\rho}_e = (\mathbb{1} \otimes X)\rho_e(\mathbb{1} \otimes X^{\dagger})$  where X is defined such that  $|\eta\rangle = \mathbb{1} \otimes X |\Omega\rangle$ .

With this we have established the missing direction in the proof of the lectures' theorem relating positive maps and observables as witnesses.

It is also possible to construct an observable that detects  $\rho_e$  itself from  $\Lambda$ . To this end, we define  $\mathcal{W}_e = (\mathbb{1} \otimes \Lambda^*)(|\eta\rangle\langle\eta|)$ .

b) Show that this construction, in fact, gives rise to an entanglement witness  $W_e$  for  $\rho_e$ .

As an application of this construction we consider the following setting. In our (fictitious) lab, we are trying to prepare a two-qubit state  $|\psi\rangle \in \mathcal{H} = \mathbb{C}^2 \otimes \mathbb{C}^2$ . We use a simple model<sup>1</sup> for what is actually happening in the lab, namely that we prepare a state with some noise

$$\rho(p) := p |\psi\rangle\langle\psi| + (1-p)\frac{1}{4}.$$

Our goal is to have an observable witness that decides whether  $\rho(p)$  is entangled or not. To this end, we will use the fact that for two-qubits system there exist no entangled PPT states. Therefore, if  $\rho(p)$  is entangled, the partial transpose  $\mathbb{1} \otimes T$  will always detect  $\rho(p)$ ..

- c) Assume  $|\psi\rangle$  has Schmidt decomposition  $|\psi\rangle = a|01\rangle + b|10\rangle$ . Determine the values of p depending on a, b such that  $\rho(p)$  is entangled.
- d) Use the eigenvector corresponding to a negative eigenvalue of  $(\mathbb{1} \otimes T)(\rho(p))$  in order to derive an entanglement witness W for  $\rho(p)$ .
- e) Show that, in fact, the witness  $\mathcal{W}$  detects all entangled states of the form  $\rho(p)$ .

Further reading: Gühne and Tóth (2009): A nice review on detecting entanglement. Ioannou et al. (2004): An algorithmic viewpoint on the problem of deciding the separability problem with a good introduction into the geometric features of entanglement.

# References

Gühne, O. and G. Tóth (2009, April). Entanglement detection. *Physics Reports* 474 (1-6), 1–75

Ioannou, L. M., B. C. Travaglione, D. Cheung, and A. K. Ekert (2004, December). Improved algorithm for quantum separability and entanglement detection. *Physical Review A* 70 (6).

<sup>&</sup>lt;sup>1</sup>What is the corresponding noise channel for this model?